

$$\beta = -\frac{1}{c} \left[ \frac{\partial \lambda}{\partial t} + c f(t) \right] + f(t)$$

$$\beta = -\frac{1}{c} \frac{\partial \lambda}{\partial t} - f(t) + f(t)$$

$$\boxed{\beta = -\frac{1}{c} \frac{\partial \lambda}{\partial t}}$$

$$\alpha = \vec{\nabla} \lambda' = \vec{\nabla} \lambda \quad \left( \because \lambda = \lambda' - c \int_0^t f(t) dt \right)$$

Therefore

$$\vec{A}' = \vec{A} + \alpha$$

$$\boxed{\vec{A}' = \vec{A} + \vec{\nabla} \lambda}$$

$$\boxed{V' = V - \frac{1}{c} \frac{\partial \lambda}{\partial t}}$$

Gauge transformation.

These are the new potential but leaves the fields  $\vec{E}'$  and  $\vec{B}$  unchanged.

Lorentz Gauge  $\Rightarrow$

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

but  $\vec{B} = \vec{\nabla} \times \vec{A}$ , and  $\vec{E} = -\vec{\nabla} V - \frac{\partial \vec{A}}{\partial t}$

$$\nabla \times (\vec{\nabla} \times \vec{A}) = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial}{\partial t} \left( -\nabla V - \frac{\partial \vec{A}}{\partial t} \right)$$

$$\begin{aligned} \therefore \nabla \times (\vec{\nabla} \times \vec{A}) &= \vec{\nabla} \cdot (\vec{\nabla} \cdot \vec{A}) - (\vec{\nabla} \cdot \vec{\nabla}) \vec{A} \\ &= \vec{\nabla} \cdot (\vec{\nabla} \cdot \vec{A}) - \nabla^2 \vec{A} \end{aligned}$$

$$\vec{\nabla} \cdot (\vec{\nabla} \cdot \vec{A}) - \nabla^2 \vec{A} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \left[ -\nabla \frac{\partial V}{\partial t} - \frac{\partial^2 \vec{A}}{\partial t^2} \right]$$

$$-\nabla^2 \vec{A} + \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} = \mu_0 \vec{J} - \vec{\nabla} \cdot (\vec{\nabla} \cdot \vec{A}) - \mu_0 \epsilon_0 \nabla \frac{\partial V}{\partial t}$$

$$-\nabla^2 \vec{A} + \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} = \mu_0 \vec{J} - \vec{\nabla} \left[ \vec{\nabla} \cdot \vec{A} + \mu_0 \epsilon_0 \frac{\partial V}{\partial t} \right]$$

$$\nabla^2 \vec{A} + \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu_0 \vec{J} + \vec{\nabla} \left[ \vec{\nabla} \cdot \vec{A} + \mu_0 \epsilon_0 \frac{\partial V}{\partial t} \right]$$

if  $\boxed{\vec{\nabla} \cdot \vec{A} + \mu_0 \epsilon_0 \frac{\partial V}{\partial t} = 0}$  Lorentz gauge.

Then  $\nabla^2 \vec{A} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu_0 \vec{J}$

Coulomb gauge  $\Rightarrow$

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\vec{E} = -\vec{\nabla} V - \frac{\partial \vec{A}}{\partial t}$$

$$\vec{\nabla} \cdot \left( -\vec{\nabla} V - \frac{\partial \vec{A}}{\partial t} \right) = + \frac{\rho}{\epsilon_0}$$

$$\nabla^2 V + \frac{\partial}{\partial t} (\vec{\nabla} \cdot \vec{A}) = - \frac{\rho}{\epsilon_0}$$

$\boxed{\vec{\nabla} \cdot \vec{A} = 0}$  Coulomb gauge.

$\nabla^2 V = - \frac{\rho}{\epsilon_0}$  (Poisson equation)

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if  $\rho = 0$ , then  $\nabla^2 V = 0$  Laplace equation.

or

$$\vec{\nabla} \cdot \vec{A}' = -\mu_0 \epsilon_0 \frac{\partial V}{\partial t}$$

(Lorentz gauge)

$$\nabla^2 V + \frac{\partial}{\partial t} \left( -\mu_0 \epsilon_0 \frac{\partial V}{\partial t} \right) = -\frac{\rho}{\epsilon_0}$$

$$\boxed{\nabla^2 V - \mu_0 \epsilon_0 \frac{\partial^2 V}{\partial t^2} = -\frac{\rho}{\epsilon_0}}$$

Boundary conditions  $\rightarrow$

Those conditions by which fields must be satisfied at the interface of the media are called boundary conditions.

To find electrostatic boundary condition we use two Maxwell equations.

$$\vec{\nabla} \cdot \vec{E} = \frac{\rho}{\epsilon_0} \Rightarrow \vec{\nabla} \cdot \vec{D} = \rho \quad (\because \vec{D} = \epsilon_0 \vec{E})$$

$$\int \vec{D} \cdot d\vec{s} = Q$$

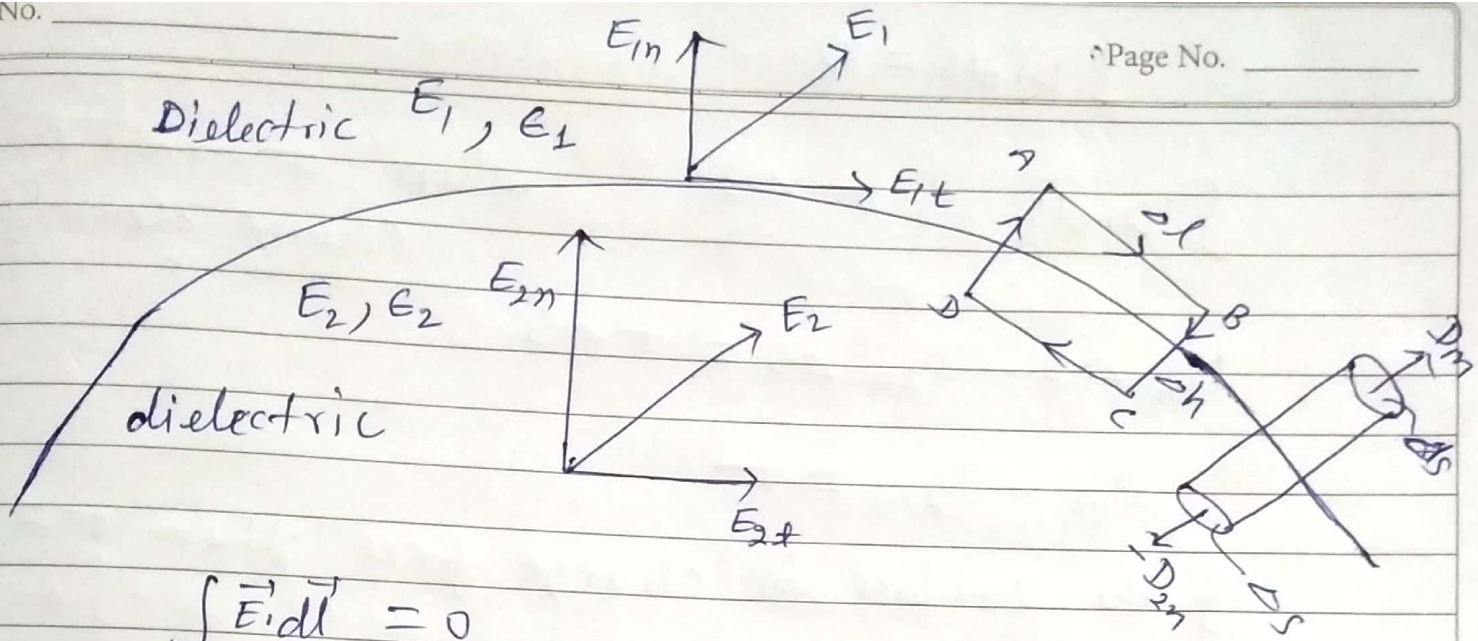
$$\vec{\nabla} \times \vec{E} = 0$$

$$\int \vec{E} \cdot d\vec{l} = 0$$

Dielectric-dielectric boundary condition  $\Rightarrow$

$$E_1 = E_{1t} + E_{1n}$$

$$E_2 = E_{2t} + E_{2n}$$



$$\int \vec{E} \cdot d\vec{l} = 0$$

$$\int_A^B \vec{E}_1 \cdot d\vec{l} + \int_B^C \vec{E}_1 \cdot d\vec{l} + \int_C^D \vec{E}_2 \cdot d\vec{l} + \int_D^A \vec{E}_2 \cdot d\vec{l} = 0$$

$$E_{1t} \Delta l + E_{1n} \frac{\Delta h}{2} - E_{2n} \frac{\Delta h}{2} - E_{2t} \Delta l + E_{2n} \frac{\Delta h}{2} + E_{1n} \frac{\Delta h}{2} = 0$$

$$(E_{1t} - E_{2t}) \Delta l = 0$$

$$\Delta l \neq 0$$

$$E_{1t} - E_{2t} = 0$$

$$\boxed{E_{1t} = E_{2t}}$$

Tangential component of  $\vec{E}$  is continuous.

$$E_{1t} = E_{2t}$$

$$\boxed{\frac{D_{1t}}{\epsilon_1} = \frac{D_{2t}}{\epsilon_2}}$$

Tangential component of  $\vec{D}$  is discontinuous across the boundary.

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$$\int \vec{D} \cdot d\vec{s} = Q$$

$$\int \vec{D} \cdot d\vec{s} = \sigma \Delta S$$

where  $\sigma$  is the surface charge density.

$$D_{1n} \Delta S - D_{2n} \Delta S = \sigma \Delta S$$

$$D_{1n} - D_{2n} = \sigma$$

if the surface is charge free then  $\sigma = 0$

$$D_{1n} - D_{2n} = 0$$

$$\boxed{D_{1n} = D_{2n}}$$

Normal component of  $D$  is continuous if the surface is charge free.

$$D_{1n} = D_{2n}$$

$$\epsilon_1 E_{1n} = \epsilon_2 E_{2n}$$

$$\boxed{\frac{E_{1n}}{\epsilon_2} = \frac{E_{2n}}{\epsilon_1}}$$

Normal component of  $\vec{E}$  is discontinuous across the boundary.